Quantum Groups and Quantizations of Isomonodromic Systems

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§1. Introduction

Isomonodromic Systems

= Isomonodromic Deformations + Discrete Symmetries

Jimbo-Miwa-Ueno, Physica 2D, 1981. Jimbo-Miwa, Physica 2D, 4D, 1981.

- Isomonodromic deformations
 - = monodromy preserving deformations (differential equations) of rational connections on $\mathbb{P}^1_{\mathbb{C}}$ (or on compact Riemann surfaces).
- Deformation parameters = time variables
 - = positions of singularities and irregular types of irregular singularities
- Discrete symmetries
 - = discrete group actions compatible with isomonodromic deformations.
 - = Bäcklund transformations of deformation differential equations

Quantizations of isomonodromic deformations

- the Schlesinger equations → the Knizhnik-Zamolodchikov equations (Reshetikhin (LMP26, 1992), Harnad (hep-th/9406078))
 The KZ equations have hypergeometric integral solutions.
- the generalized Schlesinger equations (rank-1 irreg. sing. at ∞)
- the generalized Knizhnik-Zamolodchikov equations (Babujian-Kitaev (for sl_2 , JMP39, 1998),

Felder-Markov-Tarasov-Varchenko (for any \mathfrak{g} , math.QA/0001184)) The gen. KZ equations have confluent hypergeometric integral solutions.

Conjecture. Any quantum isomonodromic system has (confluent or non-confluent) hypergeometric integral solutions.

Problem. Quantize the discrete symmetries (the Schlesinger transformations, the birational Weyl group actions, ...).

Quantizations of discrete symmetries

- the q-difference version of the birational Weyl group action (Kajiwara-Noumi-Yamada (nlin.SI/0012063))
- the quantum q-difference version of the birat. Weyl group action (Koji Hasegawa (math.QA/0703036))
- the higher Painlevé equation of type $A_l^{(1)}$ with $\widetilde{W}\big(A_l^{(1)}\big)$ symmetry (rank-2 irr. sing. at ∞) (Noumi-Yamada (math.QA/9808003))
- the quantum higher Painlevé equation type $A_l^{(1)}$ with $\widetilde{W}(A_l^{(1)})$ sym. (Hajime Nagoya (math.QA/0402281))
- the birational Weyl group action arising from a nilpotent Poisson algebra (Noumi-Yamada (math.QA/0012028))
- \downarrow complex powers of Chevalley generators in the Kac-Moody algebra the Weyl group action on the quotient skew field of $U(\mathfrak{n})\otimes U(\mathfrak{h})$

- the dressing chains (Shabat-Yamilov (LMJ2, 1991),
 (Veselov-Shabat (FAA27, 1993), V. E. Adler (Phys.D73, 1994))
 → the quantum dressing chains (Lipan-Rasinariu (hep-th/0006074))
- $\circ R(z) := z + P^{12}, \quad L_k(z) := \begin{bmatrix} x_k & 1 \\ x_k \partial_k \varepsilon_k + z & \partial_k \end{bmatrix}, \quad \partial_k = \partial/\partial x_k.$
- $R(z-w)L_k(z)^1L_k(w)^2 = L_k(w)^2L_k(z)^1R(z-w).$
- Assume n = 2g + 1, $x_{k+n} = x_k$, $\varepsilon_{k+n} = \varepsilon_k + \kappa$ (quasi-periodicity).
- The fundamental algebra of the quantum dressing chain is not the algebra generated by x_k, ∂_k but the algebra generated by $f_k := \partial_k + x_{k+1}$. The Hamiltonian of the dressing chain can be expressed with f_k .
- **Duality.** the quantum quasi-periodic dressing chain with period n \cong the quantum higher Painlevé equation of type $A_{n-1}^{(1)}$.
- \circ Thus the $\widetilde{W}(A_{2q}^{(1)})$ symmetry of the dressing chain is also quantized.

Quantizations of Isomonodromic Systems

Classical	Quantum
Poisson algebra $S(\mathfrak{g})=\mathbb{C}[\mathfrak{g}^*]$	Non-commutative algebra $U(\mathfrak{g})$
(generalized) Schlesinger eq.	(generalized) KZ eq.
$A_l^{(1)}$ higher Painlevé eq.	quantum $A_l^{(1)}$ higher Painlevé eq.
with $\widetilde{W}ig(A_l^{(1)}ig)$ symmetry	with $\widetilde{W}ig(A_l^{(1)}ig)$ symmetry
dressing chain	quantum dressing chain
with quasi-period $2g+1$	with quasi-period $2g+1$
$(\cong A_{2g}^{(1)}$ higher Painlevé eq.)	$(\cong$ quantum $A_{2g}^{(1)}$ higher Painlevé eq.)
and its $\widetilde{W}ig(A_{2g}^{(1)}ig)$ -symmetry	and its $\widetilde{W}(A_{2g}^{(1)})$ -symmetry
birational Weyl group action	the " $U_q(\mathfrak{g}) o U(\mathfrak{g})$ " limit of
arising from nilpotent Poisson	the Weyl group action on
algebra of NY	$\underline{Q(U_q(\mathfrak{n})\otimes U_q(\mathfrak{h}))}$ constructed in §2

(As far as the speaker knows, the <u>red-colored</u> results are new.)

Quantum *q*-difference Versions of Discrete Symmetries

q-difference Classical	q-difference Quantum
Poisson algebra $\mathbb{C}[G^*]$	Non-commutative algebra $U_q(\mathfrak{g})$
(G = Poisson Lie group)	(quantum universal enveloping alg.)
q-difference version of the	Weyl group action on the quotient
NY birat. Weyl group action	skew field $Q(U_q(\mathfrak{n})\otimes U_q(\mathfrak{h}))$
arising from nilp. Poisson alg.	constructed in §2
q-difference version of	quantum q -difference version of
birational Weyl Group action	birational Weyl Group action
of KNY (nlin.SI/0012063)	of Hasegawa (reconstructed in §2)
$\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$	$\underline{\operatorname{quantum}\ \widetilde{W}\big(A_{m-1}^{(1)}\big)\times\widetilde{W}\big(A_{n-1}^{(1)}\big)}$
action of KNY	action of §3

(As far as the speaker knows, the <u>red-colored</u> results are new.)

§2. Complex powers of Chevalley generators in quantum groups

Problem 1. Find a quantum q-difference version of the Noumi-Yamada birational Weyl group action arising from a nilpotent Poisson algebra (math.QA/0012028).

Answer. Using complex powers of Chevalley generators in quantum groups, we can naturally construct the quantum q-difference version of the NY birational action arising from a nilpotent Poisson algebra.

Problem 2. Find a quantum group interpretation of the quantum q-difference version of the birational Weyl group action constructed by Koji Hasegawa (math.QA/0703036).

Answer. Using complex powers of Chevalley generators in quantum groups, we can reconstruct the Hasegawa quantum birat. action.

Complex powers of Chevalley generators

- $A = [a_{ij}]_{i,j \in I}$, symmetrizable GCM. $d_i a_{ij} = d_j a_{ji}$. $q_i := q^{d_i}$.
- $U_q(\mathfrak{n}_-) = \langle f_i \mid i \in I \rangle := \text{maximal nilpotent subalgebra of } U_q(\mathfrak{g}(A)).$
- $U_q(\mathfrak{h}) = \langle a_\lambda = q^\lambda \mid \lambda \in \mathfrak{h} \rangle := \text{Cartan subalgebra of } U_q(\mathfrak{g}(A)).$
- $\alpha_i^{\vee} := \text{simple coroot}, \quad \alpha_i := \text{simple root}, \quad a_i := a_{\alpha_i} = q^{\alpha_i} = q_i^{\alpha_i^{\vee}}.$
- $\mathcal{K}_A := Q(U_q(\mathfrak{n}_-) \otimes U_q(\mathfrak{h})) = \text{the quotient skew field of } U_q(\mathfrak{n}_-) \otimes U_q(\mathfrak{h}).$
- $a_{\lambda} = q^{\lambda}$ regarded as a central element of \mathcal{K}_A is called a parameter.

Complex powers of f_i : (lohara-Malikov (hep-th/9305138))

• The action of $\mathrm{Ad}(f_i^{\lambda})x = f_i^{\lambda}xf_i^{-\lambda}$ on \mathcal{K}_A is well-defined.

$$f_i^{\lambda} f_j f_i^{-\lambda} = q_i^{\lambda} f_j + [\lambda]_{q_i} (f_i f_j - q_i^{-1} f_j f_i) f_i^{-1}$$

$$= [1 - \lambda]_{q_i} f_j + [\lambda]_{q_i} f_i f_j f_i^{-1}$$
 if $a_{ij} = -1$,

where $[x]_q := (q^x - q^{-x})/(q - q^{-1}).$

Verma relations \iff **Coxeter relations**

Verma relations of Chevalley generators f_i in $U_q(\mathfrak{n}_-)$:

$$f_i^a f_j^{a+b} f_i^b = f_j^b f_i^{a+b} f_j^a \quad (a, b \in \mathbb{Z}_{\geq 0}) \quad \text{if } a_{ij} a_{ji} = 1.$$

(formulae for non-simply-laced cases are omitted)

(Lusztig, Introduction to Quantum Groups, Prop.39.3.7 or Lemma 42.1.2.)

- Verma relations can be extended to the complex powers f_i^{λ} .
- $\tilde{r}_i \lambda \tilde{r}_i^{-1} = \lambda \langle \alpha_i^{\vee}, \lambda \rangle \alpha_i$ for $\lambda \in \mathfrak{h}$ (Weyl group action on parameters).
- Verma relations of f_i 's \iff Coxeter relations of $R_i := f_i^{\alpha_i^\vee} \tilde{r}_i$'s.

$$\circ R_i^2 = f_i^{\alpha_i^{\vee}} \tilde{r}_i f_i^{\alpha_i^{\vee}} \tilde{r}_i = f_i^{\alpha_i^{\vee}} f_i^{-\alpha_i^{\vee}} \tilde{r}_i^2 = 1.$$

$$\circ R_{i}R_{j}R_{i} = f_{i}^{\alpha_{i}^{\vee}}\tilde{r}_{i}f_{j}^{\alpha_{j}^{\vee}}\tilde{r}_{j}f_{i}^{\alpha_{i}^{\vee}}\tilde{r}_{i} = f_{i}^{\alpha_{i}^{\vee}}f_{j}^{\alpha_{i}^{\vee}+\alpha_{j}^{\vee}}f_{i}^{\alpha_{j}^{\vee}}\tilde{r}_{i}\tilde{r}_{j}\tilde{r}_{i}$$

$$= f_{j}^{\alpha_{j}^{\vee}}f_{i}^{\alpha_{i}^{\vee}+\alpha_{j}^{\vee}}f_{j}^{\alpha_{i}^{\vee}}\tilde{r}_{j}\tilde{r}_{i}\tilde{r}_{j} = f_{j}^{\alpha_{j}^{\vee}}\tilde{r}_{j}f_{i}^{\alpha_{i}^{\vee}}\tilde{r}_{i}f_{j}^{\alpha_{j}^{\vee}}\tilde{r}_{i}f_{j}^{\alpha_{j}^{\vee}}\tilde{r}_{j} = R_{j}R_{i}R_{j} \quad \text{if } a_{ij}a_{ji} = 1.$$

(formulae for non-simply-laced cases are omitted)

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Theorem. $\operatorname{Ad}(R_i) = \operatorname{Ad}(f_i^{\alpha_i^{\vee}} \tilde{r}_i)$ $(i \in I)$ generate the action of the Weyl group on \mathcal{K}_A as algebra automorphisms. This is the quantum q-difference version of the Noumi-Yamada birational Weyl group action arising from a nilpotent Poisson algebra (math.QA/0012028).

Example. If $a_{ij} = -1$, then

$$f_i^2 f_j - (q_i + q_i^{-1}) f_i f_j f_i + f_j f_i f_i = 0,$$

$$Ad(R_i)f_j = f_i^{\alpha_i^{\vee}} f_j f_i^{-\alpha_i^{\vee}} = q_i^{\alpha_i^{\vee}} f_j + [\alpha_i^{\vee}]_{q_i} (f_i f_j - q_i^{-1} f_j f_i) f_i^{-1}$$
$$= [1 - \alpha_i^{\vee}]_{q_i} f_j + [\alpha_i^{\vee}]_{q_i} f_i f_j f_i^{-1},$$

$$Ad(R_i)a_i = \tilde{r}_i a_i \tilde{r}_i^{-1} = a_i^{-1}, \quad Ad(R_i)a_j = \tilde{r}_i a_j \tilde{r}_i^{-1} = a_i a_j.$$

In particular, as the $q \rightarrow 1$ limit, we have

$$Ad(R_i)f_j = f_j + \alpha_i^{\vee}[f_i, f_j]f_i^{-1} = (1 - \alpha_i^{\vee})f_j + \alpha_i^{\vee}f_if_jf_i^{-1}.$$

Truncated q-Serre relations and Weyl group actions

Assumptions:

- $k_i k_j = k_j k_i$, $k_i f_j k_i^{-1} = q_i^{-a_{ij}} f_j$. (the action of the Cartan subalgebra)
- $f_i f_j = q_i^{\pm (-a_{ij})} f_j f_i$ $(i \neq j)$. (truncated q-Serre relations)
- $\bullet \ f_{i1}:=f_i\otimes 1$, $f_{i2}:=k_i^{-1}\otimes f_i$. $(f_{i1}+f_{i2}=\text{``coproduct of }f_i\text{''})$

Skew field \mathcal{K}_H generated by F_i, a_i :

- $\mathcal{K}_H :=$ the skew field generated by $F_i := a_i^{-1} f_{i1}^{-1} f_{i2}$, $a_i = q^{\alpha_i}$.
- Then $F_iF_j=q_i^{\pm 2(-a_{ij})}F_jF_i$ $(i \neq j)$, $a_i \in \text{center of } \mathcal{K}_H.$
- $\tilde{r}_i a_j \tilde{r}_i^{-1} = a_i^{-a_{ij}} a_j$. (the action of the Weyl group on parameters).

Theorem. Put $R_i := (f_{i1} + f_{i2})^{\alpha_i^{\vee}} \tilde{r}_i$.

Then $Ad(R_i)$'s generate the action of the Weyl group on \mathcal{K}_H .

q-binomial theorem and explicit formulae of actions

• Applying the q-binomial theorem to $f_{i1}f_{i2}=q_i^{-2}f_{i2}f_{i1}$, we obtain

$$(f_{i1} + f_{i2})^{\alpha_i^{\vee}} = \frac{(a_i^{-1} F_i)_{i,\infty}}{(a_i F_i)_{i,\infty}} f_{i1}^{\alpha_i^{\vee}}, \quad \text{where } (x)_{i,\infty} := \prod_{\mu=0}^{\infty} (1 + q_i^{2\mu} x).$$

Explicit Formulae. If $i \neq j$, then

$$Ad(R_i)F_i = F_i$$
,

$$\operatorname{Ad}(R_{i})F_{j} = \begin{cases} F_{j} \prod_{\mu=0}^{-a_{ij}-1} \frac{1 + q_{i}^{2\mu} a_{i} F_{i}}{a_{i} + q_{i}^{2\mu} F_{i}} & \text{if } f_{i}f_{j} = q_{i}^{+2(-a_{ij})} f_{j}f_{i}, \\ \prod_{\mu=0}^{-a_{ij}-1} \frac{a_{i} + q_{i}^{2\mu} F_{i}}{1 + q_{i}^{2\mu} a_{i} F_{i}} & \text{if } f_{i}f_{j} = q_{i}^{-2(-a_{ij})} f_{j}f_{i}. \end{cases}$$

• These formulae coincide with those of the quantum q-difference Weyl group action constructed by Koji Hasegawa (math.QA/0703036).

§3. Quantization of the $\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$ action of KNY

Problem 3. For any integers $m, n \geq 2$, construct

- (a) a non-commutative skew field $\mathcal{K}_{m,n}$ and
- (b) an action of $\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$ on $\mathcal{K}_{m,n}$ as alg. automorphisms which is a quantization of the Kajiwara-Noumi-Yamada action of $\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$ on $\mathbb{C}(x_{ik}|1 \leq i \leq m, 1 \leq k \leq n)$.

Answer. If m, n are mutually prime, then we can construct a quantization of the KNY action.

Tools.

- (a) Gauge invariant subalgebras of quotients of affine quantum groups,
- (b) Complex powers of corrected Chevalley generators.

The KNY discrete dynamical systems

Kajiwara-Noumi-Yamada, nlin.SI/0106029, Discrete dynamical systems with $W(A_{m-1}^{(1)} \times A_{n-1}^{(1)})$ symmetry. Kajiwara-Noumi-Yamada, nlin.SI/0112045. Noumi-Yamada, math-ph/0203030.

- (1) Action of $\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$ as algebra automorphisms on the rational function field $\mathbb{C}(x_{ik}|1 \leq i \leq m, 1 \leq k \leq n)$.
- (2) Lax representations $\implies q$ -difference isomonodromic systems.
- (3) Poisson brackets are, however, not given.

First Problem. Usually quantization replaces Poisson brackets by commutators. How to find an appropriate quantization of $\mathbb{C}(x_{ik}|1 \leq i \leq m, 1 \leq k \leq n)$ without Poisson brackets?

Minimal representations of Borel subalgebra of $U_q(\widehat{\operatorname{gl}}_m)$

• $\mathcal{B}_{m,n}$:= the associative algebra over $\mathbb{C}(q,r',s')$ generated by $a_{ik}^{\pm 1},b_{ik}^{\pm 1}$ $(i,k\in\mathbb{Z})$ with following fundamental relations:

$$a_{i+m,k} = r'a_{ik}, \quad a_{i,k+n} = s'a_{ik}, \quad b_{i+m,k} = r'b_{ik}, \quad b_{i,k+n} = s'b_{ik},$$

$$a_{ik}b_{ik} = q^{-1}b_{ik}a_{ik}, \quad a_{ik}b_{i-1,k} = qb_{i-1,k}a_{ik}.$$

All other combinations from $\{a_{ik},b_{ik}\}_{1\leq i\leq m,\ 1\leq k\leq n}$ commute.

- $U_q(\mathfrak{b}_-) = \langle t_i, f_i \mid i \in \mathbb{Z} \rangle$:= the lower Borel subalgebra of $U_q(\widehat{\mathrm{gl}}_m)$ with fundamental relations: $t_{i+m} = r't_i$, $f_{i+m} = f_i$, $t_i t_j = t_j t_i$, $t_i f_i = q^{-1} f_i t_i$, $t_i f_{i-1} = q f_{i-1} t_i$, $f_i f_j = f_j f_i$ $(j \not\equiv i \pm 1 \pmod m)$, $f_i^2 f_{i\pm 1} (q + q^{-1}) f_i f_{i\pm 1} f_i + f_{i\pm 1} f_i^2 = 0$ (q-Serre relations).
- For each k, the algebra homomorphism $U_q(\mathfrak{b}_-) \to \mathcal{B}_{m,n}$ is given by $t_i \mapsto a_{ik}$, $f_i \mapsto a_{ik}^{-1}b_{ik}$. (minimal representations of $U_q(\mathfrak{b}_-)$)

RLL = LLR relations (Quantum group)

R-matrix:
$$R(z) := \sum_{i=1}^m (q-z/q) E_{ii} \otimes E_{ii} + \sum_{i \neq j} (1-z) E_{ii} \otimes E_{jj} + \sum_{i < j} \left((q-q^{-1}) E_{ij} \otimes E_{ji} + (q-q^{-1}) z E_{ji} \otimes E_{ij} \right).$$

$$\emph{L} ext{-operators:}\quad L_k(z) := egin{bmatrix} a_{1k} & b_{1k} & & & & & \\ & a_{2k} & \ddots & & & \\ & & \ddots & b_{m-1,k} \\ b_{mk} z & & & a_{mk} \end{bmatrix}.$$

RLL = LLR relations:

$$R(z/w)L_k(z)^1L_k(w)^2=L_k(w)^2L_k(z)^1R(z/w),$$
 $L_k(z)^1L_l(w)^2=L_l(w)^2L_k(z)^1 \qquad (k\not\equiv l \ ({
m mod}\ n)),$ where $L_k(z)^1:=L_k(z)\otimes 1,\ L_k(w)^2:=1\otimes L_k(w).$

Gauge invariant subalgebra $A_{m,n} = \mathcal{B}_{m,n}^{\mathcal{G}}$ of $\mathcal{B}_{m,n}$

Gauge group:
$$\mathcal{G}:=(\mathbb{C}^{\times})^{mn}\ni g=(g_{ik}).$$
 $g_{i+m,k}=g_{ik},\ g_{i,k+n}=g_{ik}.$

Gauge transformation: The algebra automorphism of $\mathcal{B}_{m,n}$ is given by

$$a_{ik}\mapsto g_{ik}a_{ik}g_{i,k+1}^{-1}, \quad b_{ik}\mapsto g_{ik}b_{ik}g_{i+1,k+1}^{-1},$$
 i.e. $L_k(z)\mapsto g_kL_k(z)g_{k+1}^{-1} \quad (g_k:=\mathrm{diag}(g_{1k},g_{2k},\ldots,g_{mk})).$

- Assume that m, n are **mutually prime** integers ≥ 2 .
- ullet $\widetilde{m}:=\mathsf{mod}\text{-}n$ inverse of m $(\widetilde{m}m\equiv 1\ (\mathsf{mod}\ n)$, $\widetilde{m}=1,2,\ldots,n-1)$.
- The gauge invariant subalgebra $\mathcal{A}_{m,n}:=\mathcal{B}_{m,n}^{\mathcal{G}}$ of $\mathcal{B}_{m,n}$ is generated by $x_{ik}^{\pm 1}:=\left(a_{ik}(b_{ik}b_{i+1,k+1}\cdots b_{i+\widetilde{m}m-1})^{-1}\right)^{\pm 1}$, $b_{all}^{\pm 1}:=\left(\prod_{i=1}^{m}\prod_{k=1}^{n}b_{ik}\right)^{\pm 1}\in \text{center of }\mathcal{B}_{m,n}.$
- $\mathcal{K}_{m,n} :=$ the quotient skew field of $\mathcal{A}_{m,n}$ is an appropriate quantization of the rational function field $\mathbb{C}(x_{ik}|1 \leq i \leq m, 1 \leq k \leq n)$.

q-commutation relations of x_{ik} 's

- $B := \{ (\mu \mod m, \ \mu \mod n) \in \mathbb{Z}/m\mathbb{Z} \times \mathbb{Z}/n\mathbb{Z} \mid \mu = 0, 1, \dots, \widetilde{m}m 1 \}.$
- $p_{\mu\nu} := egin{cases} q & \text{if } (\mu \operatorname{mod} m, \ \nu \operatorname{mod} n) \in B, \\ 1 & \text{otherwise.} \end{cases}$
- $q_{\mu\nu} := (p_{\mu\nu}/p_{\mu-1,\nu})^2 \in \{1, q^{\pm 2}\}.$ (definition of $q_{\mu\nu}$)

Fundamental relations of x_{ik} 's:

$$x_{i+m,k} = rx_{ik}, \quad x_{i,k+n} = sx_{ik}$$
 $(r := r'^{1-\widetilde{m}m}, s := s'^{1-\widetilde{m}m}),$ $x_{i+\mu,k+\nu}x_{ik} = q_{\mu\nu}x_{ik}x_{i+\mu,k+\nu}$ $(0 \le \mu < m, 0 \le \nu < n).$

Example. If (m,n)=(2,3), then $\widetilde{m}=2$ and

$$[p_{\mu\nu}] = \begin{bmatrix} q & 1 & q \\ q & q & 1 \end{bmatrix}, \quad [q_{\mu\nu}] = \begin{bmatrix} 1 & q^{-2} & q^2 \\ 1 & q^2 & q^{-2} \end{bmatrix} \quad \begin{pmatrix} \mu = 0, 1 \\ \nu = 0, 1, 2 \end{pmatrix}.$$

$$x_{11}x_{11} = x_{11}x_{11}, \quad x_{12}x_{11} = q^{-2}x_{11}x_{12}, \quad x_{13}x_{11} = q^2x_{11}x_{13},$$

$$x_{21}x_{11} = x_{11}x_{21}, \quad x_{22}x_{11} = q^2x_{11}x_{22}, \quad x_{23}x_{11} = q^{-2}x_{11}x_{23}.$$

Example. (1) If (m,n)=(2,2g+1), then $\widetilde{m}=g+1$ and

$$[q_{\mu\nu}] = \begin{bmatrix} 1 & q^{-2} & q^2 & \cdots & q^{-2} & q^2 \\ 1 & q^2 & q^{-2} & \cdots & q^2 & q^{-2} \end{bmatrix} \quad \begin{pmatrix} \mu = 0, 1 \\ \nu = 0, 1, 2, \dots, 2g - 1, 2g \end{pmatrix}.$$

$$1 < k \le n \implies x_{1k}x_{11} = q^{(-1)^{k-1}2}x_{11}x_{1k}, \ x_{2k}x_{11} = q^{(-1)^k2}x_{11}x_{2k}.$$

(2) If (m, n) = (2g + 1, 2), then $\tilde{m} = 1$ and

$$[p_{\mu\nu}] = \begin{bmatrix} q & 1 \\ 1 & q \\ q & 1 \\ \vdots & \vdots \\ 1 & q \\ q & 1 \end{bmatrix}, \quad [q_{\mu\nu}] = \begin{bmatrix} 1 & 1 \\ q^{-2} & q^2 \\ q^2 & q^{-2} \\ \vdots & \vdots \\ q^{-2} & q^2 \\ q^2 & q^{-2} \end{bmatrix} \quad \begin{pmatrix} \mu = 0, 1, 2, \dots, 2g - 1, 2g \\ \nu = 0, 1 \end{pmatrix}.$$

Observation: $\mathcal{A}_{2,n} \cong \mathcal{A}_{n,2}$, $x_{ik} \leftrightarrow x_{ki}$, $q \leftrightarrow q$, $r \leftrightarrow s$, $s \leftrightarrow r$.

Example. (1) If (m,n)=(3,4), then $\widetilde{m}=3$ and

$$[p_{\mu\nu}] = \begin{bmatrix} q & 1 & q & q \\ q & q & 1 & q \\ q & q & q & 1 \end{bmatrix}, \quad [q_{\mu\nu}] = \begin{bmatrix} 1 & q^{-2} & 1 & q^2 \\ 1 & q^2 & q^{-2} & 1 \\ 1 & 1 & q^2 & q^{-2} \end{bmatrix} \quad \begin{pmatrix} \mu = 0, 1, 2 \\ \nu = 0, 1, 2, 3 \end{pmatrix}.$$

$$x_{12}x_{11} = q^{-2}x_{11}x_{12}$$
, $x_{13}x_{11} = x_{11}x_{13}$, $x_{14}x_{11} = q^2x_{11}x_{14}$, ...

(2) If (m,n)=(4,3), then $\widetilde{m}=1$ and

$$[p_{\mu\nu}] = \begin{bmatrix} q & 1 & 1 \\ 1 & q & 1 \\ 1 & 1 & q \\ q & 1 & 1 \end{bmatrix}, \quad [q_{\mu\nu}] = \begin{bmatrix} 1 & 1 & 1 \\ q^{-2} & q^2 & 1 \\ 1 & q^{-2} & q^2 \\ q^2 & 1 & q^{-2} \end{bmatrix} \qquad \begin{pmatrix} \mu = 0, 1, 2, 3 \\ \nu = 0, 1, 2 \end{pmatrix}.$$

Observation: $A_{3,4} \cong A_{4,3}$, $x_{ik} \leftrightarrow x_{ki}$, $q \leftrightarrow q$, $r \leftrightarrow s$, $s \leftrightarrow r$.

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Example. (1) If (m,n)=(3,5), then $\widetilde{m}=2$ and

$$[p_{\mu\nu}] = \begin{bmatrix} q & 1 & 1 & q & 1 \\ 1 & q & 1 & 1 & q \\ q & 1 & q & 1 & 1 \end{bmatrix}, \quad [q_{\mu\nu}] = \begin{bmatrix} 1 & 1 & q^{-2} & q^2 & 1 \\ q^{-2} & q^2 & 1 & q^{-2} & q^2 \\ q^2 & q^{-2} & q^2 & 1 & q^{-2} \end{bmatrix}.$$

$$x_{12}x_{11} = x_{11}x_{12}$$
, $x_{13}x_{11} = q^{-2}x_{11}x_{13}$, $x_{14}x_{11} = q^2x_{11}x_{14}$, ...

(2) If (m,n)=(5,3), then $\widetilde{m}=2$ and

$$[p_{\mu\nu}] = \begin{bmatrix} q & 1 & q \\ q & q & 1 \\ 1 & q & q \\ q & 1 & q \\ q & q & 1 \end{bmatrix}, \quad [q_{\mu\nu}] = \begin{bmatrix} 1 & q^{-2} & q^2 \\ 1 & q^2 & q^{-2} \\ q^{-2} & 1 & q^2 \\ q^2 & q^{-2} & 1 \\ 1 & q^2 & q^{-2} \end{bmatrix}.$$

Observation: $A_{3,5} \cong A_{5,3}$, $x_{ik} \leftrightarrow x_{ki}$, $q \leftrightarrow q$, $r \leftrightarrow s$, $s \leftrightarrow r$.

Symmetries of $A_{m,n}$

Duality. The algebra isomorphism $A_{m,n} \cong A_{n,m}$ is given by

$$x_{ik} \leftrightarrow x_{ki}, \quad q \leftrightarrow q, \quad r \leftrightarrow s, \quad s \leftrightarrow r.$$

Reversal. The algebra involution of $A_{m,n}$ is given by

$$x_{ik} \leftrightarrow x_{-i,-k}, \quad q \leftrightarrow q^{-1}, \quad r \leftrightarrow s^{-1}, \quad s \leftrightarrow r^{-1}.$$

Translation. For any integers μ, ν , the algebra automorphism of $\mathcal{A}_{m,n}$ is given by

$$x_{ik} \mapsto x_{i+\mu,k+\nu}, \quad q \mapsto q, \quad r \mapsto r, \quad s \mapsto s.$$

Extended affine Weyl groups $\widetilde{W}(A_{m-1}^{(1)})$, $\widetilde{W}(A_{n-1}^{(1)})$

- $W(A_{m-1}^{(1)}) := \langle r_0, r_1, \dots, r_{m-1}, \omega \rangle$ with fundamental relations: $r_i r_j = r_j r_i \ (j \not\equiv i, i+1 \ (\text{mod } m)), \ r_i r_{i+1} r_i = r_{i+1} r_i r_{i+1}, \ r_i^2 = 1, \ \omega r_i \omega^{-1} = r_{i+1} \ (r_{i+m} = r_i).$
- $T_i := r_{i-1} \cdots r_2 r_1 \omega r_{m-1} \cdots r_{i+1} r_i$ (translations).
- $\circ \widetilde{W}(A_{m-1}^{(1)}) = \langle r_1, r_2, \dots, r_{m-1} \rangle \ltimes \langle T_1, T_2, \dots, T_m \rangle \cong S_m \ltimes \mathbb{Z}^m.$
- $\widetilde{W}(A_{n-1}^{(1)}) := \langle s_0, s_1, \dots, s_{n-1}, \varpi \rangle$ with fundamental relations: $s_k s_l = s_l s_k \ (l \not\equiv k, k+1 \ (\operatorname{mod} n)), \ s_k s_{k+1} s_k = s_{k+1} s_k s_{k+1}, \ s_k^2 = 1,$ $\varpi s_k \varpi^{-1} = s_{k+1} \ (s_{k+n} = s_k).$
- $\circ U_k := s_{k-1} \cdots s_2 s_1 \varpi s_{n-1} \cdots s_{k+1} s_k$ (translations).
- $\circ \widetilde{W}(A_{n-1}^{(1)}) = \langle s_1, s_2, \dots, s_{n-1} \rangle \ltimes \langle U_1, U_2, \dots, U_n \rangle \cong S_n \ltimes \mathbb{Z}^n.$

Explicit formulae of the action of $\widetilde{W}ig(A_{m-1}^{(1)}ig)$ on $\mathcal{K}_{m,n}$

•
$$\widetilde{W}(A_{m-1}^{(1)})=\langle r_0,r_1,\ldots,r_{m-1},\omega\rangle$$
 acts on $\mathcal{K}_{m,n}=Q(\mathcal{A}_{m,n})$ by

$$r_i(x_{il}) = x_{il} - s^{-1} \frac{c_{i,l+1} - c_{i+1,l+2}}{P_{i,l+1}} = sP_{il}x_{i+1,l}P_{i,l+1}^{-1},$$
$$r_i(x_{i+1,l}) = x_{i+1,l} + s^{-1} \frac{c_{il} - c_{i+1,l+1}}{P_{i,l}} = s^{-1}P_{il}^{-1}x_{il}P_{i,l+1},$$

$$r_i(x_{il}) = x_{il} \quad (j \not\equiv i, i+1 \pmod{m}),$$

$$\omega(x_{il}) = x_{i+1,l},$$

where $c_{ik} := x_{ik}x_{i,k+1}\cdots x_{i,k+n-1}$ and

$$P_{ik} := \sum_{l=1}^{n} \underbrace{x_{ik} x_{i,k+1} \cdots x_{i,k+l-2}}_{l-1} \underbrace{x_{i+1,k+l} x_{i+1,k+l+1} \cdots x_{i+1,k+n-1}}_{n-l}.$$

Explicit formulae of the action of $\widetilde{W}ig(A_{n-1}^{(1)}ig)$ on $\mathcal{K}_{m,n}$

•
$$\widetilde{W} \left(A_{n-1}^{(1)} \right) = \langle s_0, s_1, \dots, s_{n-1}, \varpi \rangle$$
 acts on $\mathcal{K}_{m,n} = Q(\mathcal{A}_{m,n})$ by
$$s_k(x_{jk}) = x_{jk} - r^{-1} \frac{d_{j+1,k} - d_{j+2,k+1}}{Q_{j+1,k}} = rQ_{j+1,k}^{-1} x_{j,k+1} Q_{jk},$$

$$s_k(x_{j,k+1}) = x_{j,k+1} + r^{-1} \frac{d_{jk} - d_{j+1,k+1}}{Q_{jk}} = r^{-1} Q_{j+1,k} x_{jk} Q_{jk},$$

$$s_k(x_{jl}) = x_{jl} \quad (l \not\equiv k, k+1 \pmod{n}),$$

$$\varpi(x_{jl}) = x_{j,l+1},$$

where $d_{ik} := x_{i+m-1,k} \cdots x_{i+1,k} x_{ik}$ and

$$Q_{ik} := \sum_{j=1}^{m} \overbrace{x_{i+m-1,k+1} \cdots x_{i+j+1,k+1} x_{i+j,k+1}}^{m-j} \overbrace{x_{i+j-2,k} \cdots x_{i+1,k} x_{ik}}^{j-1}.$$

Duality of the extended affine Weyl group actions

- $x_{ik}^{(m,n)} := x_{ik} \in \mathcal{A}_{m,n}, \ c_{ik}^{(m,n)} := c_{ik} \in \mathcal{A}_{m,n}, \ P_{ik}^{(m,n)} := P_{ik} \in \mathcal{A}_{m,n}, \ s_i^{(m,n)} := (s_i\text{-action on } \mathcal{K}_{m,n}), \ \omega^{(m,n)} := (\omega\text{-action on } \mathcal{K}_{m,n}), \ \text{etc.}$
- The algebra isomorphism $\theta: \mathcal{A}_{m,n} \stackrel{\sim}{\to} \mathcal{A}_{n,m}$ is defined by

$$\theta(x_{ik}^{(m,n)}) = x_{-k,-i}^{(n,m)}, \ \theta(q) = q^{-1}, \ \theta(r) = s^{-1}, \ \theta(s) = r^{-1}.$$

Then

$$\theta(c_{ik}^{(m,n)}) = d_{-k-n+1,-i}^{(n,m)}, \qquad \theta(P_{ik}^{(m,n)}) = Q_{-k-n+1,-i-1}^{(n,m)}, \\ \theta(d_{ik}^{(m,n)}) = c_{-k,-i-m+1}^{(n,m)}, \qquad \theta(Q_{ik}^{(m,n)}) = P_{-k-1,-i-m+1}^{(n,m)}.$$

Therefore

$$\theta \circ r_i^{(m,n)} = s_{-i-1}^{(n,m)} \circ \theta, \quad \theta \circ \omega^{(m,n)} = (\varpi^{(n,m)})^{-1} \circ \theta,$$

$$\theta \circ s_k^{(m,n)} = r_{-k-1}^{(n,m)} \circ \theta, \quad \theta \circ \varpi^{(m,n)} = (\omega^{(n,m)})^{-1} \circ \theta.$$

Lax representations of the actions of r_i and s_k

X-operators:
$$X_{ik}=X_{ik}(z):=\begin{bmatrix}x_{ik}&1&&&&\\&x_{i+1,k}&\ddots&&&\\&&&\ddots&1&\\r^{-k}z&&&&x_{i+m-1,k}\end{bmatrix}$$
 .

(1) The action of r_i on $\{x_{1k}, \ldots, x_{mk}\}$ is uniquely characterized by

$$r_i(X_{1k}) = G_k^{(i)} X_{1k} (G_{k+1}^{(i)})^{-1}.$$

$$G_k^{(i)} := 1 + s^{-1} \frac{c_{ik} - c_{i+1,k+1}}{P_{ik}} E_{i+1,i} \quad (c_{ik} = x_{ik} x_{i+1,k} \cdots x_{i+m-1,k}),$$

$$G_k^{(0)} := 1 + r^{k-1} z^{-1} s^{-1} \frac{c_{mk} - c_{m+1,k+1}}{P_{mk}} E_{1m}. \quad (E_{ij}'s \text{ are matrix units.})$$

$$G_k^{(0)} := 1 + r^{k-1} z^{-1} s^{-1} rac{c_{mk} - c_{m+1,k+1}}{P_{mk}} E_{1m}$$
. $(E_{ij}$'s are matrix units.)

(2) The action of s_k is uniquely characterized by

$$s_k(X_{ik}X_{i,k+1}) = X_{ik}X_{i,k+1}, \ s_k(X_{il}) = X_{il} \ (l \not\equiv k \pmod{n}),$$

 $s_k: d_{ik} \leftrightarrow d_{i+1,k+1} \ (d_{ik} = x_{i+m-1,k} \cdots x_{i+1,k}x_{ik}).$

Quantum q-difference isomonodromic systems

Monodromy matrix: $X_{ik}(z) := X_{ik}(z) X_{i,k+1}(z) \cdots X_{i,k+n-1}(z)$.

Matrix q-difference shift operator (shift parameter = s):

$$T_{z,s}v(s) := diag(s^{-1}, s^{-2}, \dots, s^{-m})v(s^m z)$$
 ($v(z)$ is m -vector valued).

Linear q-difference equation: $T_{z,s}v(z) = X_{11}(z)v(z)$.

Connection matrix preserving transformations:

- (1) $s_k(X_{11}(z)) = X_{11}(z)$ for k = 1, 2, ..., n 1.
- (2) $\varpi(X_{11}(z)) = X_{11}^{-1}X_{11}(z)X_{1,n+1} = T_{z,s}X_{1,n+1}^{-1}T_{z,s}^{-1}X_{11}(z)X_{1,n+1}.$
- $U_k = s_{k-1} \cdots s_2 s_1 \varpi s_{n-1} \cdots s_{k+1} s_k$.

The action of $\langle U_1, U_2, \dots, U_n \rangle \cong \mathbb{Z}^n$

- \longrightarrow Quantum q-difference isomonodromic dynamical system with n time variables
- ullet The action of $\widetilde{W}\big(A_{m-1}^{(1)}\big)$ \longrightarrow Symmetry of the dynamical system

Gen KUROKI (Tohoku Univ.)

Example $((m,n)=(3,2)) \bullet x_{i+3,k}=rx_{ik}, x_{i,k+2}=sx_{ik}.$

- $x_{11}x_{11} = x_{11}x_{11}$, $x_{21}x_{11} = q^{-2}x_{11}x_{21}$, $x_{31}x_{11} = q^2x_{11}x_{31}$, $x_{12}x_{11} = x_{11}x_{12}$, $x_{22}x_{11} = q^2x_{11}x_{22}$, $x_{32}x_{11} = q^{-2}x_{11}x_{32}$.
- $P_{ik} = x_{i+1,k+1} + x_{ik}$, $Q_{ik} = x_{i+2,k+1}x_{i+1,k+1} + x_{i+2,k+1}x_{ik} + x_{i+1,k}x_{ik}$.
- $r_1(x_{11}) = s(x_{22} + x_{11})x_{21}(x_{13} + x_{12})^{-1}$, $r_1(x_{21}) = s^{-1}(x_{22} + x_{11})^{-1}x_{21}(x_{13} + x_{12})$, $\omega(x_{ik}) = x_{i+1,k}$.
- $s_1(x_{11}) = r(x_{42}x_{32} + x_{42}x_{21} + x_{31}x_{21})^{-1}x_{12}(x_{32}x_{22} + x_{32}x_{11} + x_{21}x_{11}),$ $s_1(x_{12}) = r^{-1}(x_{42}x_{32} + x_{42}x_{21} + x_{31}x_{21})x_{11}(x_{32}x_{22} + x_{32}x_{11} + x_{21}x_{11})^{-1},$ $\varpi(x_{ik}) = x_{i,k+1}.$ $(U_1 = \varpi r_1, U_2 = r_1\varpi)$ $U_1(x_{11}) = r(x_{43}x_{33} + x_{43}x_{22} + x_{32}x_{22})^{-1}x_{13}(x_{33}x_{23} + x_{33}x_{12} + x_{22}x_{12}).$
- U_1 generates quantum qP_{IV} (q-difference Panlevé IV system). The action of $\widetilde{W}(A_2^{(1)})$ is symmetry of quantum qP_{IV} .

Gen KUROKI (Tohoku Univ.)

Action of
$$\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$$
 on $\mathcal{K}_{m,n}$ as alg. autom.

Theorem. For any mutually prime integers $m, n \geq 2$, the action of $\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$ on $\mathcal{K}_{m,n} = Q(\mathcal{A}_{m,n})$ as algebra automorphisms is constructed. This is a quantization of the KNY action of $\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$ on $\mathbb{C}(x_{ik}|1 \leq i \leq m, 1 \leq k \leq n)$.

Easy Part. Lax representations \Longrightarrow braid relations of r_i and s_k .

Difficult Part. To show that

 r_i and s_k act on $\mathcal{K}_{m,n} = Q(\mathcal{A}_{m,n})$ as algebra automorphisms.

Sketch of proof. Let φ_i be appropriately corrected Chevalley generators in $\mathcal{B}_{m,n}$ and put $\rho_i := \varphi_i^{\alpha_i^\vee} \tilde{r}_i$. Then $\mathrm{Ad}(\rho_i) x_{jl} = \rho_i x_{jl} \rho_i^{-1} = r_i(x_{jl})$. Therefore r_i acts on $\mathcal{K}_{m,n}$ as algebra automorphisms. The duality leads to that s_k also acts on $\mathcal{K}_{m,n}$ as algebra automorphisms.

Chevalley generators F_i

Monodromy matrix: $\mathbb{L}(z) := L_1(z)L_2(z)\cdots L_n(z)$.

($\mathbb{L}(z)$ is the product of the L-operators of the minimal representations.)

$$\mathbb{L}(z) = \begin{bmatrix} A_1 & B_1 & \cdots & \ddots & \vdots \\ & A_2 & \cdots & \ddots & \vdots \\ & & \ddots & B_{m-1} \\ 0 & & & A_m \end{bmatrix} + z \begin{bmatrix} \ddots & \ddots & \ddots & \ddots & \vdots \\ \ddots & \ddots & \ddots & \ddots & \vdots \\ B_m & \ddots & \ddots & \ddots & \vdots \end{bmatrix} + \cdots$$

- $R(z/w)\mathbb{L}(z)^1\mathbb{L}(w)^2 = \mathbb{L}(w)^2\mathbb{L}(z)^1R(z/w)$ $\implies F_i := A_i^{-1}B_i$ satisfy the q-Serre relations.
- $R_i := F_i^{\alpha_i^{\vee}} \tilde{r}_i$ generate the Weyl group action on the skew field generated by A_i , B_i , and parameters $a_{\lambda} = q^{\lambda}$ ($\lambda \in \mathfrak{h}$).
- But the action of $Ad(R_i)$ does not preserve the skew field generated by $x_{ik} = a_{ik}(b_{ik}b_{i+1,k+1}\cdots b_{i+\tilde{m}m-1})^{-1}$, b_{all} , and parameters $a_{\lambda} = q^{\lambda}$.

Correction factors for F_i

- $\mathcal{K}_{m,n} = Q(\mathcal{A}_{m,n}) \subset Q(\mathcal{B}_{m,n}).$
- $x \simeq y \iff \exists c \in (\text{the center of } Q(\mathcal{B}_{m,n}))^{\times} \text{ s.t. } cx = y.$
- $\widetilde{n} := \mathsf{mod}\text{-}m$ inverse of n ($\widetilde{n}n \equiv 1 \pmod{m}$), $\widetilde{n} = 1, 2, \ldots, m-1$).
- $v_{ik}:=b_{ik}b_{i+1,k+1}\cdots b_{i+\widetilde{n}n-1,k+\widetilde{n}n-1}$ (v_{1k} are correction factors). (cf. $x_{ik}=a_{ik}(b_{ik}b_{i+1,k+1}\cdots b_{i+\widetilde{m}m-1})^{-1}$, $\widetilde{m}=\mathsf{mod}\text{-}n$ inverse of m)
- $c_{i1}^{-1}P_{i1} \simeq v_{i1}^{-1}F_i = v_{i1}^{-1}A_i^{-1}B_i$ (motivation to find v_{ik}).
- $\varphi_i := v_{i1} F_i = v_{i1} A_i^{-1} B_i \simeq v_{i1}^2 c_{i1}^{-1} P_{i1}$ (corrected F_i).
- Using φ_i instead of F_i , we can construct the action of the affine Weyl group $W(A_{m-1}^{(1)})$ on $\mathcal{K}_{m,n}=Q(A_{m,n})$ as algebra automorphisms.

Generators of the $W(A_{m-1}^{(1)})$ -action on $\mathcal{K}_{m,n}=Q(\mathcal{A}_{m,n})$

- $\mathcal{H}_m := \mathbb{C}(q, r', s')[q^{\pm 2\varepsilon_i^{\vee}}]_{i=1}^m$, $\varepsilon_i^{\vee} := E_{ii} \in \mathfrak{h}$, $\alpha_i^{\vee} := \varepsilon_i^{\vee} \varepsilon_{i+1}^{\vee}$.
- ullet $\mathcal{A}_{m,n}\cong (\mathcal{A}_{m,n}\otimes \mathcal{H}_m)/I$ $(\otimes =\otimes_{\mathbb{C}(q,r',s')})$,

where I := the two-sided ideal generated by $c_{ii} \otimes 1 - 1 \otimes q^{-2\varepsilon_i^{\vee}}$.

- $\tilde{r}_i \varepsilon_i^{\vee} \tilde{r}_i^{-1} = \varepsilon_{i+1}^{\vee}$, $\tilde{r}_i \varepsilon_{i+1}^{\vee} \tilde{r}_i^{-1} = \varepsilon_i^{\vee}$, $\tilde{r}_i \varepsilon_j^{\vee} \tilde{r}_i^{-1} = \varepsilon_j^{\vee}$ $(j \neq i, i+1)$.
- $\rho_i := \varphi_i^{\alpha_i^{\vee}} \tilde{r}_i$. (generators of the $W(A_{m-1}^{(1)})$ -action on $\mathcal{K}_{m,n}$)
- $\operatorname{Ad}(\rho_i)$ generate the action of $W(A_{m-1}^{(1)})$ on $Q(\mathcal{A}_{m,n}\otimes\mathcal{H}_m)$.
- The actions of $\mathrm{Ad}(\rho_i)$'s on $Q(\mathcal{A}_{m,n}\otimes\mathcal{H}_m)$ induce the actions of $r_i\in W(A_{m-1}^{(1)})$ on $\mathcal{K}_{m,n}=Q(\mathcal{A}_{m,n})$:

$$Ad(\rho_i)x_{il} = r_i(x_{il}) = sP_{il}x_{i+1,l}P_{i,l+1}^{-1},$$

$$Ad(\rho_i)x_{i+1,l} = r_i(x_{i+1,l}) = s^{-1}P_{il}^{-1}x_{il}P_{i,l+1},$$

$$Ad(\rho_i)x_{jl} = r_i(x_{jl}) = x_{jl} \quad (j \not\equiv i, i+1 \pmod{m}).$$

Summary of Results

- §2. (for any symmetrizable GCM $A = [a_{ij}]$)
- Ad-action of complex powers of Chevalley generators f_i in $U_q(\mathfrak{g})$
- \implies the action of the Weyl group on $Q(U_q(\mathfrak{n}) \otimes U_q(\mathfrak{h}))$ (quantum q-difference version of the NY math.QA/0012028 action)
- ⇒ Reconstruction of the Hasegawa math.QA/0703036 action
- §3. (for any mutually prime integers $m, n \ge 2$)
- $\mathcal{B}_{m,n}:=$ the minimal representation of $U_q(\mathfrak{b})^{\otimes n}\subset U_q(\widehat{\mathrm{gl}}_m)^{\otimes n}$.
- $\mathcal{K}_{m,n} := Q$ (the gauge invariant subalgebra $\mathcal{A}_{m,n}$ of $\mathcal{B}_{m,n}$)
- $\implies \mathcal{K}_{m,n} = \text{Quantization of } \mathbb{C}(x_{ik}|1 \leq i \leq m, 1 \leq k \leq n).$
- ullet Complex powers of the **corrected** Chevalley generators in $\mathcal{B}_{m,n}$
- $\Longrightarrow \widetilde{W}(A_{m-1}^{(1)})$ -action on $\mathcal{K}_{m,n}$
- $\Longrightarrow \widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$ -action on $\mathcal{K}_{m,n}$ (by the $m \leftrightarrow n$ duality)

Other Problems

Problem. Construct commuting Hamiltonians in $U_q(\mathfrak{n}) \otimes U_q(\mathfrak{h})$ with Weyl group symmetry.

Hint. Commuting transfer matrices for " $AL^1BL^2 = CL^2DL^1$ " algebras. $(F = q^{-\sum H_i \otimes H^i}, A = P(F)^{-1}RF, B = F, C = P(F), D = R)$

Problem. Construct commuting Hamiltonians in $\mathcal{A}_{m,n}$ with $\widetilde{W}(A_{m-1}^{(1)}) \times \widetilde{W}(A_{n-1}^{(1)})$ symmetry.

Classical Case. $\det(\mathbb{X}_{11}^{(m,n)}(z)-(-1)^nw)=\det(\mathbb{X}_{11}^{(n,m)}(w)-(-1)^mz)$ generates the invariants of birational $\widetilde{W}(A_{m-1}^{(1)})\times\widetilde{W}(A_{n-1}^{(1)})$ action.

Problem. Construct solutions of quantum (q-)isomonodromic systems.

Conjecture. Schrödinger equation of any quantum (q-)isomonodromic system has (non-confluent or confluent) (q-)hypergeometric solutions.